# 46415 Structural Analysis and Design Optimization of Wind Turbine Blades

# Mini-Project 2: Optimum Blade

# (2 page report is due on the 13th of June 2018 at 13.00)

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# 1. Introduction

Mini-Project 2 builds on Mini-Project 1 introducing the following changes and new concepts:

- A 3D beam finite element model is now used.
- The thin-walled ellipse is added to the cross section as a simple representation of the airfoil.
- The wall thickness e of the ellipse is added as design variable.
- A edgewise bending loadcase is added.
- Eigenfrequency constraints are added.

# 2. Optimization Problem

The considered blade design optimization now reads

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Figure 1: Cross section of section i.



where  $m(\mathbf{a}, \mathbf{t}, \mathbf{e})$  is the mass of the blade and  $\mathbf{a} = [a_1, \dots, a_N]^T$ ,  $\mathbf{t} = [t_1, \dots, t_N]^T$  and  $\mathbf{e} = [e_1, \dots, e_N]^T$  are the design variables.  $a_i$  is the cap width at section i,  $t_i$  is the cap thickness at section i and  $e_i$  is the thickness of the ellipse at section i, as shown in Figure 1. The tip displacement  $\delta(\mathbf{a}, \mathbf{t}, \mathbf{e})$ , the bending strain  $\varepsilon_i(a_i, t_i, e_i)$  in each section and the the buckling coefficient  $\eta_i(a_i, t_i)$  in each section must be smaller or equal to given values. The eigenfrequencies of the blade are constrained by an upper and lower limit.  $N$  is the number of sections and  $k$  is the number of eigenfrequency constraints.

## 3. Analysis

The theory underlying the functions used for the analysis of the structural response of the beam are presented in here. The subsections mimic the same structure adopted in the Matlab implementation, i.e., analysis of cross section properties in compute\_csprops, assembly of the cross section constitutive matrices in compute\_constitutive, assemble of the beam finite element model in compute\_beam, and solve the beam finite element equations in compute\_beam\_solution.

#### 3.1. Assumptions and Parameters

Is is assumed that the caps of the box girder and the ellipse representing the airfoil are made from unidirectional plies and triaxial plies, respectively. The effective material properties of the unidirectional and triaxial plies are given in Table 1. It may again be assumed that the plies are available in any thickness, so the cap thickness  $t_i$  and the thickness of the ellipse  $e_i$  can be treated as a continuous variables.



Figure 2: Thinwalled ellipse for definition of cross section properties.

The height  $h(z)$ , the chord length  $c(z)$ , and the bending moments  $M_x(z)$  and  $M_y(z)$  are described by given spline curves, which can be loaded into Matlab using the command load('splines'). The splines are also defined in Appendix A. The first derivative of the spline curves describing bending moments yields the shear force, the second derivative yields the distributed load (force per length).

The box girder is partitioned into  $N = 20$  sections of equal length  $\Delta z$ . Table 1 lists a number of model parameters.

## 3.2. Cross Section Stiffness and Mass Properties

The cross section stiffness and mass properties are discussed in this section. See the function compute\_csprops for a Matlab implementation of the theory presented here.

## 3.2.1. Area

The total area is the sum of the area of the caps and ellipse. The evaluation of the area of the caps has been presented in Mini-project 1. The area of the thin-walled ellipse with constant wall thickness  $t$  shown if Figure 2 is [1, Table A.1/26]:

$$
A_{ellipse} = \pi t (a+b) \left[ 1 + K_1 \left( \frac{a-b}{a+b} \right)^2 \right] \quad \text{where} \tag{1}
$$

$$
K_1 = 0.2464 + 0.002222 \left(\frac{a}{b} + \frac{b}{a}\right) \quad . \tag{2}
$$

Note that the definition of  $a$  and  $t$  in Figure 2 differs from the global definition in Figure 1. The total area is given by

$$
A = A_{caps} + A_{ellipse} \tag{3}
$$

#### 3.2.2. Mass per unit length

The total mass per unit length  $m_s$  is the sum of the contributions from the caps and ellipse and is thus given as

$$
m_s = m_{s,caps} + m_{s,ellipse} = \rho_{uniax} A_{caps} + \rho_{triax} A_{ellipse}
$$
\n<sup>(4)</sup>



Table 1: Model parameters.

### 3.2.3. Moment of inertia

The second moment of area  $I_x$  of the thin-walled ellipse with constant wall thickness t shown if Figure 2 is  $[1,$  Table A.1/26]:

$$
I_{x, ellipse} = \frac{\pi}{4}ta^2(a+3b)\left[1+K_2\left(\frac{a-b}{a+b}\right)^2\right] + \frac{\pi}{16}t^3(3a+b)\left[1+K_3\left(\frac{a-b}{a+b}\right)^2\right] \tag{5}
$$

$$
K_2 = 0.1349 + 0.1279 \frac{a}{b} - 0.01284 \left(\frac{a}{b}\right)^2 \tag{6}
$$

$$
K_3 = 0.1349 + 0.1279 \frac{b}{a} - 0.01284 \left(\frac{b}{a}\right)^2 \tag{7}
$$

For  $I_{y,ellipse}$  interchange a and b in the expressions for  $I_{x,ellipse}$ ,  $K_2$  and  $K_3$ .

The second mass moment of inertia is obtained as

$$
I_{xx} = \rho I_x \text{ and } I_{yy} = \rho I_y. \tag{8}
$$

Consequently, the total second mass moment of inertia is given as

$$
I_{xx} = I_{xx,caps} + I_{xx,ellipse} = \rho_{uniax} I_{x,caps} + \rho_{triax} I_{x,ellipse}
$$
\n(9)

$$
I_{yy} = I_{yy,caps} + I_{yy,ellipse} = \rho_{uniax} I_{y,caps} + \rho_{triaz} I_{y,ellipse}
$$
\n
$$
(10)
$$

#### 3.2.4. Longitudinal Stiffness

The longitudinal stiffness of the beam is henceforth referred to as EA. It is obtained as the sum of the longitudinal stiffness of the caps and ellipse and is given as

$$
EA = E_{1,uniax} A_{caps} + E_{1,triax} A_{ellipse}
$$
\n<sup>(11)</sup>

where  $A_{ellipse}$  is defined in (1).

## 3.2.5. Shear Stiffness

As a simplification both shear correction factors of the ellipse are assumed to be constant and identical:

$$
k_{ellipse} = k_{x,ellipse} = k_{y,ellipse} = 0.53
$$
\n<sup>(12)</sup>

The shear stiffness of the ellipse is:

$$
k_x G A_{ellipse} = k_y G A_{ellipse} = k_{ellipse} G_{12, triax} A_{ellipse}
$$
\n
$$
(13)
$$

As the shear webs of the box girder are not included in the model, it is not possible to compute a realistic shear stiffness for the box girder. Therefore, a very high value is assigned to the shear stiffness of the caps in  $x$ - and  $y$ -direction, in fact turning the Timoshenko beam model into a Euler-Bernoulli beam model.

Note that the shear stiffness of a cross section cannot be computed as the sum of the shear stiffnesses of parts of the cross section. Formally, the respective cross section stiffness properties of the caps and the ellipse are added in the Matlab code to compute the "total" cross section stiffness properites. In the case of the shear stiffness this does not matter, because of the assumptions regarding the shear stiffness of the caps described above.

#### 3.2.6. Bending Stiffness

The bending stiffness of the beam is the sum of the bending stiffness of the caps and the bending stiffness of the ellipse.

The bending stiffness around the x and y axis of the cross section is henceforth referred to as  $EI_x$  and  $EI_y$ , respectively. It is obtained as the sum of the bending stiffness of the caps and ellipse and is given as

$$
EI_x = E_{1,uniax}I_{x,caps} + E_{1,triax}I_{x,ellipse}
$$
\n
$$
(14)
$$

$$
EI_y = E_{1,uniax}I_{y,caps} + E_{1,triax}I_{y,ellipse}
$$
\n
$$
(15)
$$

## 3.2.7. Torsional Stiffness

The torsional stiffness constant  $K$  of the thin-walled ellipse with constant wall thickness t shown if Figure 2 is  $[1,$  Table 10.1/13]:

$$
K_{ellipse} = \frac{4\pi^2 t \left[ \left( a - \frac{t}{2} \right)^2 \left( b - \frac{t}{2} \right)^2 \right]}{U} \tag{16}
$$

$$
U = \pi(a+b-t) \left[ 1 + 0.258 \frac{(a-b)^2}{(a+b-t)^2} \right]
$$
 (17)

As a simplification the contribution of the caps to the torsional stiffness of the beam is neglected and thus the torsional stiffness GK is given by

$$
GK = G_{12, triax} K_{ellipse} \tag{18}
$$

Note that the torsional stiffness of a cross section (like the shear stiffness) cannot be computed as the sum of the torsional stiffnesses of parts of the cross section. Formally, the respective cross section stiffness properties of the caps and the ellipse are added in the Matlab code to compute the "total" cross section stiffness properites. In the case of the torsional stiffness this does not matter, because the torsional stiffness of the caps is set to zero in the Matlab code.

## 3.3. Cross Section Stiffness and Mass Matrix

It is assumed that the load application point, the elastic center, the mass center, and the position of the beam finite element node coincide at each cross section. Moreover, all relevant cross section properties are determined with respect to this position.

For a linear elastic beam there exists a linear relation between the cross section generalized forces **T** and moments **M** in  $\boldsymbol{\theta} = [\mathbf{T}^T \mathbf{M}^T]^T$ , and the resulting strains  $\boldsymbol{\tau}$  and curvatures  $\kappa$  in  $\psi = \left[\tau^T \kappa^T\right]^T$ . This relation is given in its stiffness form as  $\mathbf{K}_s \psi = \boldsymbol{\theta}$ , where  $\mathbf{K}_s$ is the  $6 \times 6$  cross section stiffness matrix. In the most general case, considering material anisotropy and inhomogeneity, all the 21 stiffness parameters in  $\mathbf{K}_s$  may be required to describe the deformation of the cross section. In the current project, the entries of  $\mathbf{K}_s$  are determined as

$$
\mathbf{K}_s = \begin{bmatrix} k_x G A & 0 & 0 & 0 & 0 & 0 \\ 0 & k_y G A & 0 & 0 & 0 & 0 \\ 0 & 0 & E A & 0 & 0 & 0 \\ 0 & 0 & 0 & E I_x & 0 & 0 \\ 0 & 0 & 0 & 0 & E I_y & 0 \\ 0 & 0 & 0 & 0 & 0 & G K \end{bmatrix}
$$
(19)

The 6  $\times$  6 cross section mass matrix  $\mathbf{M}_s$  relates the linear and angular velocities in  $\phi$ to the generalized inertial linear and angular momentum in  $\gamma$  through  $\phi = M_s \gamma$ . The coefficients of  $M_s$  for the general case are

$$
\mathbf{M}_s = \begin{bmatrix} m_s & 0 & 0 & 0 & 0 & -m_s y_m \\ 0 & m_s & 0 & 0 & 0 & m_s x_m \\ 0 & 0 & m_s & m_s y_m & -m_s x_m & 0 \\ 0 & 0 & m_s y_m & I_{xx} & -I_{xy} & 0 \\ 0 & 0 & -m_s x_m & -I_{xy} & I_{yy} & 0 \\ -m_s y_m & m_s x_m & 0 & 0 & 0 & I_{xx} + I_{yy} \end{bmatrix}
$$
(20)

where  $m_s$  is the mass per unit length,  $I_{xx}$  and  $I_{yy}$  are the mass moment of inertia with respect to  $x$  and  $y$ , respectively, and  $I_{xy}$  is the product of inertia. The off-diagonal terms are due to the offset between the position of the cross section reference center and the mass center  $\mathbf{m}_c = (x_m, y_m)$ . In this project the  $\mathbf{m}_c = (0, 0, 0, 0, 0)$  and thus the cross section mass matrix is given as

$$
\mathbf{M}_s = \begin{bmatrix} m_s & 0 & 0 & 0 & 0 & 0 \\ 0 & m_s & 0 & 0 & 0 & 0 \\ 0 & 0 & m_s & 0 & 0 & 0 \\ 0 & 0 & 0 & I_{xx} & 0 & 0 \\ 0 & 0 & 0 & 0 & I_{yy} & 0 \\ 0 & 0 & 0 & 0 & 0 & I_{xx} + I_{yy} \end{bmatrix}
$$
(21)

The theory presented in this section is implemented in compute\_constitutive in Matlab.

### 3.4. Beam Finite Element Analysis

#### 3.4.1. Beam Finite Element Stiffness and Mass Matrix

The beam finite element stiffness and mass matrix for element b are given by

$$
\mathbf{K}_b = \int_0^{L_b} \mathbf{B}_b^T \mathbf{K}_s \mathbf{B}_b \, \mathrm{d}z \quad \text{and} \quad \mathbf{M}_b = \int_0^{L_b} \mathbf{N}_b^T \mathbf{M}_s \mathbf{N}_b \, \mathrm{d}z \tag{22}
$$

where  $L_b$  is the length of element b. The beam finite element stiffness matrix  $\mathbf{K}_b$  for element b is given in function of  $\mathbf{B}_b = \mathcal{B}(\mathbf{N}_b)$  where  $\mathcal{B}$  is the strain-displacement relation which is a function of  $N_b$ , the finite element shape function matrix. The cross section stiffness and mass matrices  $\mathbf{K}_s$  and  $\mathbf{M}_s$ , respectively, have been defined in the previous section. The global beam stiffness and mass matrix  $\bf{K}$  and  $\bf{M}$  are defined as

$$
\mathbf{K} = \sum_{b=1}^{n_b} \mathbf{K}_b \text{ and } \mathbf{M} = \sum_{b=1}^{n_b} \mathbf{M}_b
$$
 (23)

where  $n_b$  is the number of elements in the beam finite element assemblage, and  $\mathbf{K}_b$  and  $M_b$  are the beam finite element stiffness and mass matrix for element b, respectively. The summation refers to the typical finite element assembly. The cross section stiffness and mass matrix,  $\mathbf{K}_s$  and  $\mathbf{M}_s$ , are defined in (19) and (21), respectively.

The theory presented in this section is implemented in compute\_beam in Matlab.

#### 3.4.2. Displacement Solution

The deformation  $\bf{u}$  resulting from the loads  $\bf{f}$  is the solution to the following linear system of equations

$$
Ku = f \tag{24}
$$

where  $\bf{K}$  is the beam finite element stiffness matrix defined in (23).

Implementation notes: In Matlab the solution to a system like the one above is obtained by writing u=K\f. The theory presented in this section is implemented in compute\_beam\_solution.

#### 3.4.3. Cross Section Forces and Moments

The cross section forces  $T$  and moments  $M$  are determined based on the displacement solution obtained from (24). For the four node beam finite element from FRANS, the forces and moments at each element of the beam finite element assembly  $f_e = [\mathbf{T}_{e,1} \mathbf{M}_{e,1} \dots \mathbf{T}_{e,4} \mathbf{M}_{e,4}]$ are determined as

$$
\mathbf{f}_e = \mathbf{K}_e \mathbf{u}_e \tag{25}
$$

where  $\mathbf{u}_e$  are the entries of the displacement vector **u** which are associated with the element e, and  $\mathbf{T}_{e,k}$  and  $\mathbf{M}_{e,k}$  are the cross section forces and moments, respectively, at node k of element e. The vector of cross section strains and curvatures are obtained using the cross section constitutive relation in its compliance form, i.e.,

$$
\Psi = \mathbf{K}_s^{-1} \mathbf{f}_e \tag{26}
$$

The theory presented in this section is implemented in compute\_beam\_solution.

#### 3.4.4. Bending Strains and Buckling Analysis

The bending strains and buckling coefficient are analyzed using the same routines as in Mini-project 1. The only difference is that the curvatures in  $\kappa$  are determined from the beam finite element assembly as detailed in Section 3.4.3. Also, for buckling calculations, the thickness of the caps and ellipse are considered.

Implementation notes: Note that the cross section forces and moments are analyzed at the node closest to the root of the blade as this is the node at which the bending moment is higher. As a result the strains measured using the beam finite element are slightly higher than those measured in Mini-project 1 in which the cross section forces and moments were measured at the center of the element.

## 3.4.5. Eigenfrequency solution

The finite element form of the beam structural eigenvalue problem is

$$
\left(\mathbf{K} - \omega_f^2 \mathbf{M}\right) \mathbf{v}_f = 0, \ \forall f = 1, ..., n_d \tag{27}
$$

where  $n_d$  is the number of degrees of freedom associated with the finite element stiffness and mass matrices,  $K$  and  $M$ , respectively. The problem above yields the eigenfrequencies  $\boldsymbol{\omega} = {\{\omega_1, ..., \omega_{n_d}\}}$  associated with the eigenvectors  $\mathbf{v} = {\{\mathbf{v}_1, ..., \mathbf{v}_{n_d}\}}$ .

**Implementation notes:** In the current implementation the eigenfrequencies in  $\omega$  are given in ascending order of magnitude, i.e.,  $\omega_1 \leq \omega_2 \leq ... \leq \omega_{n_d}$ . The eigenvectors are ordered accordingly and are already mass normalized. The theory presented in this section is implemented in compute\_beam\_solution.

# 4. Sensitivity analysis

# 4.1. Cross Section Stiffness and Mass Matrix

The gradients of the cross section stiffness matrix are given by

$$
\frac{\partial \mathbf{K}_s(\mathbf{x})}{\partial x_i} = \begin{bmatrix}\n\frac{\partial k_x G A(\mathbf{x})}{\partial x_i} & 0 & 0 & 0 & 0 & 0 \\
0 & \frac{\partial k_y G A(\mathbf{x})}{\partial x_i} & 0 & 0 & 0 & 0 \\
0 & 0 & \frac{\partial E A(\mathbf{x})}{\partial x_i} & 0 & 0 & 0 \\
0 & 0 & 0 & \frac{\partial E I_x(\mathbf{x})}{\partial x_i} & 0 & 0 \\
0 & 0 & 0 & 0 & \frac{\partial E I_y(\mathbf{x})}{\partial x_i} & 0 \\
0 & 0 & 0 & 0 & \frac{\partial E I_y(\mathbf{x})}{\partial x_i} & 0 \\
0 & 0 & 0 & 0 & \frac{\partial G K(\mathbf{x})}{\partial x_i}\n\end{bmatrix}
$$
\n(28)

which resolves to the calculation of the gradients of each of the entries. The gradient of the cross section mass matrix is given by

$$
\frac{\partial \mathbf{M}_s(\mathbf{x})}{\partial x_i} = \begin{bmatrix}\n\frac{\partial m_s(\mathbf{x})}{\partial x_i} & 0 & 0 & 0 & 0 & 0 \\
0 & \frac{\partial m_s(\mathbf{x})}{\partial x_i} & 0 & 0 & 0 & 0 \\
0 & 0 & \frac{\partial m_s(\mathbf{x})}{\partial x_i} & 0 & 0 & 0 \\
0 & 0 & 0 & \frac{\partial I_{xx}(\mathbf{x})}{\partial x_i} & 0 & 0 \\
0 & 0 & 0 & 0 & \frac{\partial I_{yy}(\mathbf{x})}{\partial x_i} & 0 \\
0 & 0 & 0 & 0 & 0 & \frac{\partial I_{yy}(\mathbf{x})}{\partial x_i} + \frac{\partial I_{yy}(\mathbf{x})}{\partial x_i}\n\end{bmatrix}
$$
\n(29)

Implementation Notes: The gradients of each of the entries has been derived in the Maple file accompanying the code.

#### 4.2. Finite Element Stiffness and Mass Matrix

The sensitivities of the global beam finite element stiffness matrix  $\bf{K}$  are obtained through differentiation of  $\bf{K}$  in (23) to yield

$$
\frac{\partial \mathbf{K}(\mathbf{x})}{\partial x_i} = \sum_{b=1}^{n_b} \int_0^{L_b} \mathbf{B}_b^T \frac{\partial \mathbf{K}_s(\mathbf{x})}{\partial x_i} \mathbf{B}_b \, \, \mathrm{d}z \tag{30}
$$

The gradient of the cross section stiffness matrix  $\mathbf{K}_s$  is described in Section 4.1. The gradients of the global beam finite element mass matrix  $M(x)$  are obtained through differentiation of  $M$  in  $(23)$  and defined as

$$
\frac{\partial \mathbf{M}(\mathbf{x})}{\partial x_i} = \sum_{b=1}^{n_b} \int_0^{L_b} \mathbf{N}_b^T \frac{\partial \mathbf{M}_s(\mathbf{x})}{\partial x_i} \mathbf{N}_b \, dz \tag{31}
$$

Implementation Notes: Note that the same routines used for computing the element stiffness matrices K and M can be used to build the gradients by simply providing  $\frac{\partial \mathbf{K}_s(\mathbf{x})}{\partial x_i}$ and  $\frac{\partial \mathbf{M}_s(\mathbf{x})}{\partial x_i}$  instead of  $\mathbf{K}_s$  and  $\mathbf{M}_s$ .

#### 4.2.1. Displacement

The solution to the linear system of equations in  $(24)$  yields the displacements **u**. The gradients of the displacement with respect to the design variables for the case of design independent loads is given by

$$
\mathbf{K}\frac{\partial \mathbf{u}}{\partial \mathbf{x}} = -\frac{\partial \mathbf{K}}{\partial \mathbf{x}}\mathbf{u}
$$
 (32)

which is obtained after applying the chain rule to (24).

**Implementation Notes:** Similarly to the displacement solution in (24), the gradients  $\frac{\partial u}{\partial x}$ are obtained as dudx = - K \ (dKdx\*u). Note that this corresponding to solving the system for as many right hand sides as number of design variables. The procedure can be implemented efficiently by considering elementwise computations of the right hand side.

#### 4.2.2. Eigenfrequencies

The solution to the structural eigenvalue problem in (27) yields the eigenfrequencies and eigenvectors  $\boldsymbol{\omega} = {\{\omega_1, ..., \omega_{n_d}\}}$  and  $\mathbf{v} = {\{\mathbf{v}_1, ..., \mathbf{v}_{n_d}\}}$ , respectively. It is assumed that the eigenvectors are mass-normalized such that

$$
\mathbf{v}_p^T \mathbf{M}(\mathbf{x}) \mathbf{v}_q = \delta_{pq}, \ \forall p, q = 1, ..., n_d.
$$

where  $n_d$  is the number of degrees of freedom, and  $\delta_{pq}$  is the Kronecker delta such that  $\delta_{pq} = 1$  if  $p = q$  and  $\delta_{pq} = 0$  otherwise. The gradient of a single eigenfrequency  $\omega_p$  with respect to the design variable  $x_i$  is given by

$$
\frac{\partial \omega_p^2(\mathbf{x})}{\partial x_i} = \mathbf{v}_p^T \left( \frac{\partial \mathbf{K}(\mathbf{x})}{\partial x_i} - \omega_p^2(\mathbf{x}) \frac{\partial \mathbf{M}(\mathbf{x})}{\partial x_i} \right) \mathbf{v}_p
$$
\n(33)

Implementation Notes: Note that in order to solve the eigenvalue problem above the boundary conditions on  $K$  and  $M$  are enforced by removing the rows and columns corresponding to the d.o.f. where the beam is clamped. This is different from the approach employed for enforcing the boundary conditions when solving for the displacements.

### 5. Implementation

The focus of this section is on the description of the data necessary for running the optimization procedure, and the output from the structural analysis functions. A series of auxiliary symbols and parameters are described in Table 2.



Table 2: Symbols used for describing the fields in Tables 3 and 4

# 5.1. Input

The input stored in the structure problem, required to run the optimization, is described in Table 3.







Table 3: Description of all the fields included in the structure problem containing the parameters necessary to solve the optimization problem in Matlab.

## 5.2. Analysis Functions

The analysis of the structural response of the blade is based on the code FRANS. This is a beam finite element code developed at DTU Wind Energy for the analysis of straight beams whose section properties are described by a 6x6 cross section stiffness and mass matrix. FRANS is a linear elastic analysis tool based on four node beam elements with cubic Lagrangian shape functions.

The analysis of the total mass of the blade in compute\_objective which is used as the objective function is relatively simple. Most of the work is carried out for the calculation of the constraints in the function compute\_constraints. Here there are four functions which compose the analysis part, namely:

- [ csprops ] = compute\_csprops( problem, design ) Function for evaluation of the cross section stiffness mass and stiffness properties and its gradients.
- [ constitutive ] = compute\_constitutive( problem, csprops ) Function to assemble the cross section constitutive stiffness and mass matrix and its gradients.
- [ beam ] = compute\_beam( problem, constitutive ) Function to assemble the beam finite element stiffness and mass matrix and calculate the gradients of the element stiffness and mass matrices.

• [ solution ] = compute\_beam\_solution( problem, beam, constitutive ) - Function to calculate the solution to the beam finite element analysis displacement and eigenvalue problems.

The content of each of the structures is described in detail in the next section. These results are the building blocks for the calculation of the gradients of the constraint functions.

# 5.2.1. Output of Analysis Functions

The output from the analysis of the structural response of the beam is summarized in Table 4.





Table 4: Description of all the fields included in the structures output by the analysis functions in Matlab. The information contained in these structures are the building blocks for the implementation of the gradients.

# 6. Mini-project tasks

In order to complete the mini-project please perform (at least) the following tasks. A summary of the most important results must be included in the report.

- Download the Matlab program for solving the optimization problem from DTU Inside.
- Study the description of the project and the Matlab code (basis for the final project).
- A small part of the program is missing and needs to be completed: The computation of the gradients of the eigenfrequencies in analysis/compute\_frequency.m
- Check your implementation of the gradients of the eigenfrequencies using the finite difference checks in fmincon. If necessary implement your own finite difference checker of the user supplied gradients.
- Perform a sensitivity analysis on the lower and upper bounds on the frequency constraints by solving a number of different problems. What happens to the optimal design? What happens to the optimal mass?

• Perform the sensitivity analysis on the design driving constraints using the Lagrange multipliers. Using the Lagrange multipliers, discuss what constraints and variable bounds would have the greatest impact on the final design. Discuss in terms of relative (*i.e.* varying constraints by  $\pm 1\%$ ) and absolute variations (*i.e.* varying constraints by  $\pm 0.001$  units) what constraints are driving the design. Which constraint is the objective most sensitive to? Pick the most sensitive constraint and change the constraint bound by 10%, solve the optimization problem. Comment on whether the objective (and design) improved as you expected it. How did your lagrange multipliers change? How can this information be used in an engineering design process.

# References

[1] Warren C. Young and Richard G. Budynas. Roark's Formulas for Stress and Strain. McGraw-Hill, 2002.

## Appendix A Splines describing geometrical properties and loading

Splines are functions defined piecewise by polynomials (see Figure 3). At the intersections of two polynomial pieces, continuity conditions (e.g. n-times continuously differentiable) are usually assigned in order to achieve a smooth curve.

A spline  $p(x)$  can be described in terms of its breaks  $\xi_1, \xi_2, \ldots, \xi_l$  and its polynomial coefficients  $c_{ii}$ :

$$
p_j(x) = \sum_{i=1}^k (x - \xi_j)^{k-i} c_{ji} \qquad j = 1, 2, \dots l \quad , \tag{34}
$$

where  $l$  is the number of polynomial pieces and  $k$  is the number of coefficients in each polynomial ( $k = 4$  for a cubic spline). The polynomial  $p_i(x)$  describes the spline in the interval  $\xi_j \leq x \leq \xi_{j+1}$ .

In this project splines haven been used to describe bending moments  $(M_x(z), M_y(z))$  and geometrical properties  $(h(z), c(z))$  as a function of the radial coordinate z.

Tables 5 to 8 display the breaks and polynomial coefficients of the splines used in this project.

## Putting together splines in Matlab

The example below demonstrates how a spline defined by its breaks and coefficients can be put together in Matlab using the ppmak command.

```
breaks = [2.8000 4.8000 18.8310 27.1510 37.4240 63.6150 89.1660];
coefs = [
0.000 4.800 0.0000E+00 0.0000E+00 0.0000E+00 0.0000E+00 5.3800E+00;
4.800 18.000 7.1197E-04 -2.0874E-02 0.0000E+00 5.3800E+00;
18.000 35.000 -1.3286E-04 7.3198E-03 -1.7892E-01 3.3804E+00;35.000 80.000 -3.2478E-06 5.4401E-04 -4.5231E-02 1.8015E+00;
80.000 89.166 -3.9380E-04 8.3219E-04 -1.6000E-02 5.7179E-01];
rel_thick_spline = ppmak ( breaks , coefs ,1);
```


Figure 3: Cubic spline approximating data.

		$\xi_{j+1}$	$c_{i1}$	$c_{i2}$	$c_{i3}$	$c_{i4}$
$\vec{\mathbb{I}}$	0.000	4.800	$0.0000E + 00$	$0.0000E + 00$	$0.0000E + 00$	5.3800E+00
$j=2$	4.800	18.000	7.1197E-04	$-2.0874E-02$	$0.0000E + 00$	5.3800E+00
j=3	18.000	35.000	$-1.3286E - 04$	7.3198E-03	$-1.7892E-01$	3.3804E+00
$j=4$	35.000	80.000	3.2478E-06	5.4401E-04	$-4.5231E-02$	1.8015E+00
Cπ.	80.000	89.166	$-3.9380E - 04$	8.3219E-04	$-1.6000E-02$	5.7179E-01

Table 5: Spline describing the profile height  $h(z)$  in m as a function of the z-coordinate in m.

	5	$\xi_{j+1}$	$c_{j1}$	$c_{i2}$	$c_{i3}$	$c_{j4}$
$\equiv$	0.000	8.196	$0.000000 + 00$	$0.0000E + 00$	$0.000000 + 00$	5.3800E+00
$j=2$	8.196	19.955	$-4.6130E - 04$	1.0423E-02		5.3800E+00
$j=3$	19.955	28.012	9.5406E-05	$-5.8502E-03$	5.3769E-02	$6.0711E + 00$
$j=4$	28.012	38.222	7.9431E-05	$-3.5440E-03$	$-2.1926E - 02$	$6.1745E + 00$
ç=i	38.222	55.027	2.4247E-05	$-1.1111E-03$	$-6.9452E-02$	5.6657E+00
j≡6	55.027	70.058	8.5976E-06	1.1132E-04	$-8.6253E - 02$	4.2999E+00
$\leq \equiv 1$	70.058	78.159	4.7121E-06	4.9899E-04	$-7.7080E - 02$	$3.0578E + 00$
$j=8$	78.159	85.000	$-4.3908E - 04$	6.1351E-04	$-6.8067E-02$	$2.4686E + 00$
j=9	85.000	86.252	$-1.3287E-03$	$-8.3983E-03$	$-1.2133E-01$	1.8911E+00
$j=10$	86.252	88.659	$-1.4760E - 02$	$-1.3389E - 02$	$-1.4861E-01$	$1.7234E + 00$
$j=11$	88.659	88.986	$-6.7982E + 00$	$-1.1999E-01$	$-4.6970E - 01$	$1.0821E + 00$
$j=12$	88.986	89.166	1.0436E+01	$-6.7809E + 00$	$-2.7235E + 00$	6.7906E-01

Table 6: Spline describing the chord length  $c(z)$  in m as a function of the z-coordinate in m.

		$\xi_i$ $\xi_{i+1}$	$c_{i1}$	$C_{12}$	C <sub>13</sub>		Ci5	်ဂ်
		$0.000$ 70.000	$-6.0100E-04$	5.8905E-02	$3.0765E + 01$	2.5734E-14	$-6.8673E + 05$	$3.9149E + 07$
	70.000		80.000   -1.7972E-02	$-1.5145E-01$	1.7810E+01	$6.1311E + 03$	$-2.2581E + 05$	$2.0344E + 06$
$j=3$	80.000	87.000	$-2.5632E - 01$	$-1.0501E + 00$	$-6.2202E + 00$	$6.3948E + 03$	$-9.9352E + 04$	$4.0393E + 05$
	$j=4$   87.000		89.166 $\vert$ -3.6869E+01 -1.0021E+01		$-1.6122E + 02$		$5.0763E + 03$ $-1.5257E + 04$	1.2846E+04

Table 7: Spline describing the bending moment  $M_x(z)$  in Nm as a function of the zcoordinate in m.



Table 8: Spline describing the bending moment  $M_y(z)$  in Nm as a function of the zcoordinate in m.